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## 8. Molecular dynamics simulations on tensile and torsional deformation and strength of super-nanotubes

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**Abstract.** In this book chapter the mechanical properties of single walled carbon nanotubes (SWCNT) arranged in a form of super-tubes (STs), under both tensile and torsion, are investigated using classical molecular dynamics simulations, based on reactive empirical bond-order potential. The investigated STs presented values of fracture toughness, fracture energy and dissipated energy that are only 5, 8, and 2 times smaller than the ones presented by the constituent SWCNT, respectively, but are much more flexible. Based on the predicted mechanical properties, STs may represent new candidates for novel porous, flexible and high strength and tough materials, e.g. as scaffolds in the regenerative medicine.

### 1. Introduction

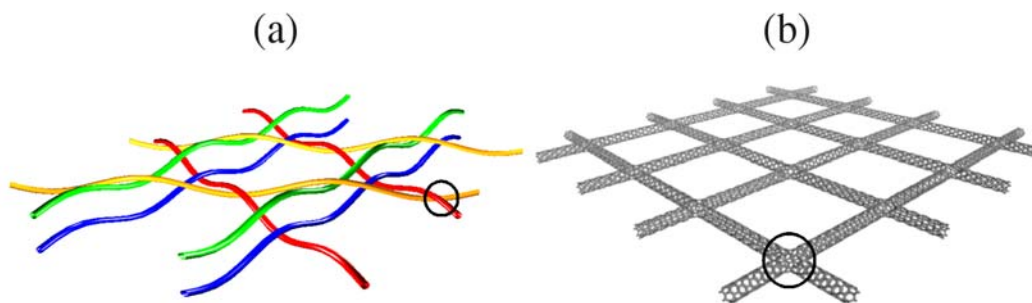
Carbon nanotubes, possessing extremely high dynamic strength and stiffness, are ideal candidates as nanofiber-reinforcements [1]. Many attempts have been made to develop procedures to controllably assemble large number

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of single walled carbon nanotubes (SWCNTs) in terms of position and orientation. These procedures would allow the fabrication of ordered SWCNT networks and their use in designing of new materials with desirable electronic and mechanical properties. Many research groups have been working on the production of carbon nanotube networks [2-8]. Snow *et al.* have demonstrated the capability of random carbon nanotube networks to function as transistors [3] and chemical sensors [4]. Many companies such as DuPont, IBM, Intel, Motorola, and Samsung are currently doing research in order to develop applications based on carbon nanotube networks [9]. Recently, Hall *et al.* showed that carbon nanotube networks can exhibit unusual mechanical properties such as negative Poisson's ratio, i.e., a lateral dimension expands during stretching [10].

With the purpose of producing *ordered* carbon nanotube networks, Diehl *et al.* have developed a technique that applies electric field to controllably arrange nanotubes on a “crossbar” configuration on silicon oxide surfaces [7]. Exploring the morphology of silicon carbide surfaces, Derycke *et al.* were able to produce ordered nanotube networks with “hexagonal” and “crossbar” arrangements [2]. These nanotubes were produced after the heating of the surface over 1500 °C in vacuum. Scanning tunneling microscopy images indicated that the nanotube organization follows the surface morphology. This technique of using surfaces has been applied by other groups to produce highly aligned carbon nanotubes on the surface [11,12] and networks composed exclusively by ordered SWCNTs [12]. In those networks, the nanotubes are weakly bonded by van der Waals interactions with the nanotubes radially deformed at the junctions (Fig. 1 (a)).

The ultimate arrangement involving SWCNTs would be the one where isolated SWCNTs are covalently connected by using X-, Y-, or T-like junctions,



**Figure 1.** Schematic view of (a) a random and (b) an ordered carbon nanotube network. In this example, the elements of the random network are weakly bonded at the junctions (indicated by the black circles) while the ones of the ordered mat are covalently connected.

forming an highly-ordered network (Fig. 1 (b)). The mechanical properties of these ideal networks were investigated by using molecular dynamics simulations [13]. The simulations indicated that these ordered networks become stiffer under large displacements. This is associated to the high Young's modulus presented by the nanotubes ( $\sim 1$  TPa) and due to the difference of the effective spring constants related to the nanotube bending at the junctions and to the nanotube axial strain.

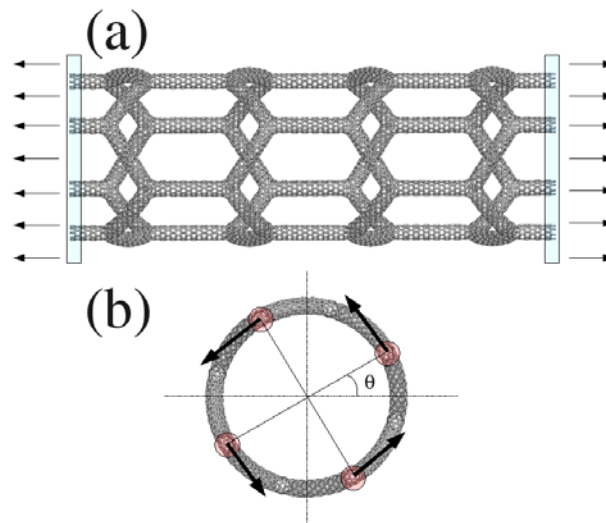
In this work we investigate the rupture of ordered single walled carbon nanotube networks using atomistic simulations. The rupture process is analyzed in networks formed by rolling up planar mats during tensile and torsional deformation regimes. These rolled up mats were named “super” carbon nanotubes (STs) [14] and, similarly to a  $(n,m)$  SWCNT [1],  $(N,M)$  ST with different chiralities can be constructed. The nomenclature  $[N,M]@(n,m)$  indicates a ST of chirality  $(N,M)$  constructed with  $(n,m)$  SWNTs.

## 2. Methodology

The rupture and mechanical properties of ordered SWCNT networks have been investigated using molecular dynamics simulations [15]. The adaptive intermolecular reactive empirical bond-order (AIREBO) potential developed by Stuart *et al.* [16] was used to model carbon-carbon interactions. AIREBO is similar to the reactive potential developed by Brenner [17] but it incorporates by suitable modifications the non-bonded interactions through an adaptative treatment of the intermolecular interactions. This kind of reactive potentials have been proved to be accurate to describe carbon nanotube deformations under mechanical strain [18]. During the molecular dynamics simulations, Newton's equations of motion were integrated with a third-order Nordieck predictor corrector algorithm [19] using a time step of 0.5 fs.

We have carried out simulations of tensile tests in situations of impact load, i.e., the atoms on the ST extremities were moved along the axial directions with a speed of 10 m/s (Fig. 2). Torsional strains were applied through the rotation of the atoms on the ST extremities (radius fixed) in a rate of  $1^\circ$  per 1000 time steps ( $3.4 \cdot 10^{-5}$  rad/fs).

In order to investigate thermal effects on the tensile behaviour of STs, the Berendsen thermostat [20] was applied to all remaining atoms. Tensile and torsional behaviours for 300 K were analysed. We have investigated zigzag  $[N, 0]@(8, 0)$  and armchair  $[M, M]@(8, 0)$  STs with  $4 \leq N \leq 6$  and  $4 \leq M \leq 6$ , with initial tube lengths ranging from 8.6 to 35.2 nm, corresponding to two ST unit cells.

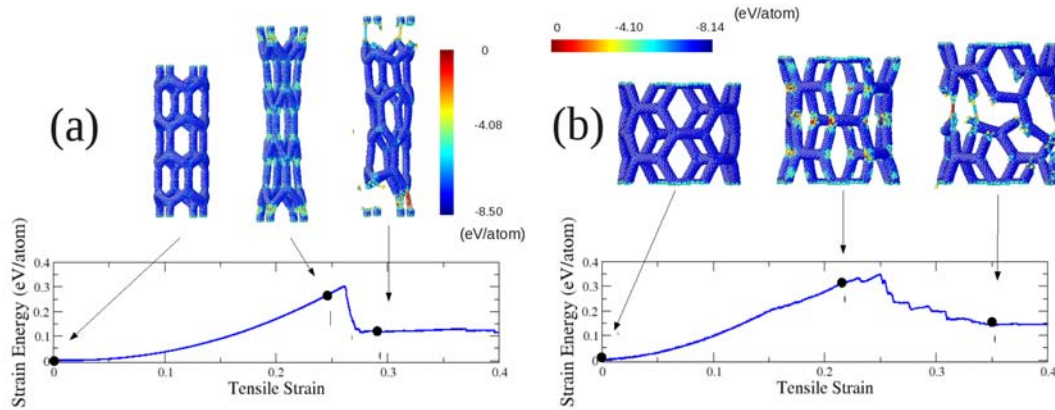


**Figure 2.** A carbon nanotube network (“super” carbon nanotube) under (a) tensile and (b) torsional strains. Longitudinal (a) and cross section (b) views of a  $[4, 0]@(8, 0)$  ST. In the ST construction single-walled carbon nanotubes are connected by Y-like junctions. The atoms in the ST extremities (filled boxes, a) are moved with a constant speed during tensile tests. Filled discs (b) represent the cross-sectional ST area.

### 3. Results

Figure 3 shows the tensile behavior for a zig-zag and an armchair ST. The rupture in zigzag STs occurs at breaking strains at about 28–30% (Fig. 3 (a)). This value is close to the one obtained for a  $(8, 0)$  SWCNT. The ST deformation is mainly due to angle changing rather than SWCNT stretching, as confirmed by the observed hyperelasticity due to a fishing net like behaviour. During the tensile deformation the angles between SWCNTs in the ST structures are changed and the stress concentration is mainly observed at the junctions. The ST rupture occurs near the junctions close to the ST extremities where different SWCNTs break approximately at the same time, causing an abrupt decrease of the strain energy. Thus, due to the large changes in the angles between SWCNTs, the rupture happens before the occurrence of an effective stretching of SWCNTs. In fact, on increasing the zigzag ST radius the breaking strain is shifted to a higher value than the one found in a zigzag ST with a smaller radius.

For the armchair ST the rupture occurs not abruptly, as in the zigzag case, but in small steps through a multistep breaking of the junctions. This behaviour can be seen in the strain energy evolution shown in Fig. 3 (b), where many decreasing steps of the strain energy are observed and associated with local ruptures on the ST structure.



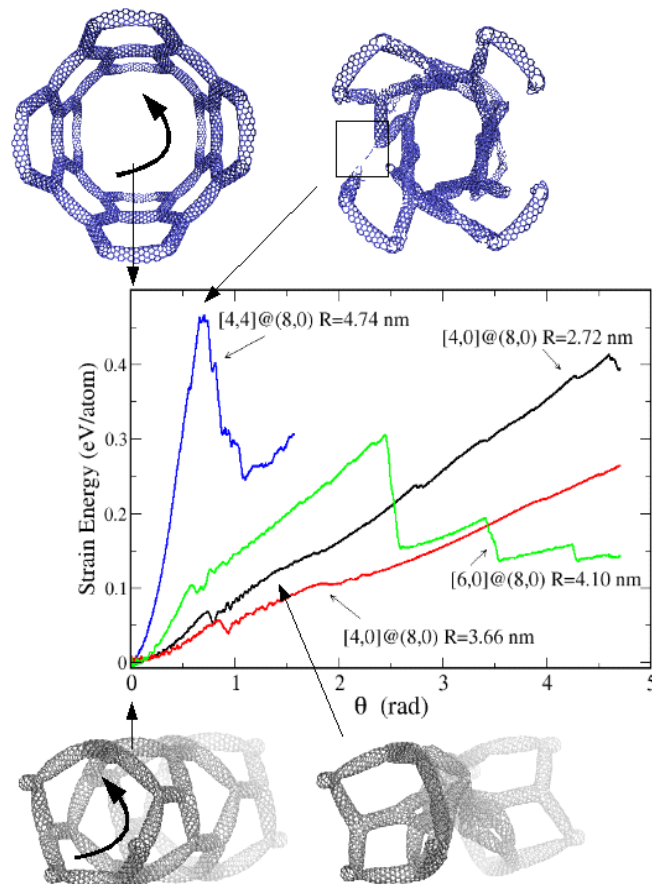
**Figure 3.** Behaviour of the strain energy as function of the tensile strain for the (a)  $[4, 0]@(8, 0)$  ST with  $R = 2.72$  nm and (b)  $[4, 4]@(8, 0)$  ST with  $R = 4.74$  nm at 300 K. Snapshots of the simulation are presented for specific tensile strains. Atoms are coloured according to their potential energy.

The key points related to the different behaviours presented by zigzag and armchair STs are the junctions and their arrangement in the STs. We have associated the deformation behaviour of STs with two main mechanisms. The first one is associated with the SWCNT stretching, dictated by the SWCNT mechanical property. The second one is due to the deformation at the junctions, i.e., changes in the angle of the junctions. These changes are relatively easier to occur than the SWCNT stretching due to the high value of the SWCNT Young's modulus. Consequently, for cases where the arrangement of junctions in the ST facilitates the deformation of the first type in a tensile test, the ST will present a tensile strength and breaking strain similar to the constituent SWCNT. It will not be exactly the same due to the effect of the second type of deformation. This is the case for the  $[4, 0]@(8, 0)$  STs (Fig. 3 (a)). On the other hand, when the second type of deformation is favoured during the tensile test due to the arrangement of the junctions, the angle changes become more relevant than the SWCNT stretching, as in the case of the  $[4, 4]@(8, 0)$  STs (Fig. 3 (b)).

Table 1 summarizes some of the mechanical properties of zig-zag and armchair STs investigated here. We can use these data to estimate the toughness presented by STs. Using the estimated reduction ( $\sim 50\%$ ) of dissipated energy per unit mass presented by STs compared to (8, 0) SWCNTs at 300 K and the value of carbon nanotube fibres comprising SWCNTs in a polymer matrix (570 J/g [21]), that is higher than that presented by spider dragline silk (165 J/g [22]) and Kevlar (33 J/g [22]), we can roughly estimate a real toughness value for STs as  $\sim 570/2 = 285$  J/g.

**Table 1.** Atomistic results for the tensile strength  $\sigma_c$ , breaking strain  $\varepsilon_c$ , fracture toughness ( $K_c$ ), fracture energy ( $G_c$ ), and dissipated energy per unit mass ( $D_c$ ) for 300 K. The values correspond to a single simulation [24]. The MD predicted values of the fracture toughness, fracture energy and of the dissipated energy per unit mass of the STs must be considered realistic only if compared with those of the SWCNT.

Structure	Radius (nm)	$\sigma_c$ (GPa)	$\varepsilon_c$	$K_c$ (MPa m <sup>1/2</sup> )	$G_c$ (N/m)	$D_c$ (kJ/g)
(8, 0)	0.31	92.6	0.28	25.6	876	10.7
[4, 0]@(8, 0)	2.72	79.2	0.26	6.8	160	5.5
[4, 0]@(8, 0)	3.66	77.9	0.27	5.8	188	4.9
[4, 4]@(8, 0)	4.74	55.5	0.22	6.2	106	6.5
[4, 4]@(8, 0)	6.36	36.5	0.15	4.5	81	3.6



**Figure 4.** Behaviour of the torsional strain energy as function of the rotation angle  $\theta$  for different STs for 300 K. Snapshots of the simulation are presented for specific  $\theta$  values for the [4,4]@(8,0) (top) and [4,0]@(8,0) with radius  $R=2.72$  nm (bottom). The square on the top left snapshot shows the region of a junction that broke at about an angle of  $\theta=45^\circ$ .

Therefore, our calculations, resulting in strengths of several gigapascals and toughness of  $\sim 280$  J/g, suggest that super-composites based on STs [23] could be competitive with super-tough carbon nanotube fibres (strength of 1.8 GPa and toughness of 570 J/g [21]).

The torsional behaviour of some STs is shown in Figure 4. From our simulations we noticed that zigzag STs can achieve larger torsional strain before breaking compared to armchair STs. For the  $[4,0]@(8,0)$  ST with a radius of 2.72 nm we observed a collapse of the structure when  $\theta$  is about  $45^\circ$ , which is associated to the first decrease on the strain energy. The collapsed structure is similar to the one shown on the bottom left of Fig. 4. As the radius increases the angle where collapse occurs is shifted to large values. Similarly to the case of tensile strains, the rupture for torsional strains occurs mainly on the regions of the nanotube junctions as we can see for the  $[4,4]@(8,0)$  ST in Fig. 4. The breaking of the tube due to torsional strain occurs at about  $40^\circ$  and  $137^\circ$  for  $[4,4]@(8,0)$  and  $[6,0]@(8,0)$  ( $R=4.10$  nm) STs, respectively. On the other hand, for the  $[4,0]@(8,0)$  STs the rupture was observed for values above  $250^\circ$ . These results indicate that STs also present high flexibility under torsional loads. We obtained the following values for the second derivative of the torsional strain energy with respect to the rotation angle (in eV/atom/rad<sup>2</sup>): 0.12, 0.08, 0.13, and 1.10, for the  $[4,0]@(8,0)$  ( $R=2.72$  nm),  $[4,0]@(8,0)$  ( $R=3.66$  nm),  $[6,0]@(8,0)$  ( $R=4.10$  nm), and  $[4,4]@(8,0)$  ( $R=4.74$  nm), respectively.

## 4. Conclusions

The mechanical properties of single walled carbon nanotubes arranged in a form of super-tubes are investigated using classical molecular dynamics simulations based on reactive empirical bond-order potential. From tensile tests of impact loads, we have found that STs are more flexible than the SWCNT used to form them but, in some cases, they show comparable tensile strengths. During tensile deformations the shape and aperture of pores in ST sidewalls can be modified providing a way to vary the accessible channels to the inner parts of STs through the application of mechanical loads. Torsional tests indicated ST chirality dependence on the breaking torsional strain values. The ST rupture occurs basically at regions near the SWCNT junctions and it is influenced by the ST chirality. The investigated STs presented values of fracture toughness, fracture energy, and dissipated energy that are about 5, 8, and 2 times smaller than the ones presented by the constituent (8,0) SWCNT, respectively. Simulations indicated that these networks are also very flexible under torsional loads, mainly zigzag STs. Based on the

predicted mechanical properties, STs may represent new candidates for novel porous, flexible, and high strength materials.

## 5. Acknowledgements

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